

1 **Surface layer formation in the earliest stages of corrosion of steel in CO<sub>2</sub>-saturated brine**  
2 **at 80°C, studied by *in situ* synchrotron X-ray methods**

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12

13 **Abstract**

14 Grazing-incidence small-angle X-ray scattering (GISAXS) from polycrystalline steel shows  
15 features associated with the underlying microstructure, and, in the initial stages of  
16 corrosion, with the development of very small-scale surface roughness, on the nanometer  
17 height scale. A 1Cr0.25Mo pipeline steel in hot, CO<sub>2</sub>-saturated brine develops the very  
18 small-scale surface roughness significantly faster than a simple carbon steel, although the  
19 overall dissolution current density for the two steels is almost the same. We speculate that

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20 it is due to the presence of a layer comprising 'blobs' of amorphous  $\text{FeCO}_3$ , which grow to  
21 spread over the surface and eventually cover it, because it is significantly larger in height  
22 scale than the roughness developed during the initial stages of anodic dissolution in  
23 acidified NaCl solution, where no surface film is expected. The greater roughness on the  
24 1Cr0.25Mo steel can be interpreted as due to small pre-crystalline nuclei, that form at much  
25 lower supersaturation, and grow faster, than on mild steel. Grazing-incidence X-ray  
26 diffraction studies show at later stages the apparent preferential dissolution of smaller  
27 crystallites of iron, with spatial size scale 0.1 – 1  $\mu\text{m}$ . This develops significantly more slowly  
28 on the 1Cr0.25Mo than on the simple carbon steel.

29 Keywords: small-angle X-ray scattering, anodic film, carbon dioxide, corrosion, iron  
30 carbonate

### 31 **Introduction**

32 The formation of crystalline layers of siderite is important for the diminution of corrosion  
33 rate of carbon steel in hot,  $\text{CO}_2$ -saturated brine – a phenomenon of some importance for  
34 management of corrosion of oil and gas production pipelines (1). The corrosion reaction is  
35 complex (2): iron dissolves as a carbonato complex of some sort (3), to form a colloidal  
36 product assumed to be amorphous iron carbonate (4); the growth of siderite requires a  
37 significant critical supersaturation (4-6); and the supersaturation is controlled by the  
38 transport of dissolved  $\text{CO}_2$ , the dissolution rate of iron and the precipitation rate of the  
39 colloidal iron carbonate (2). Dissolution of iron is presumed to occur through a surface layer  
40 of amorphous iron carbonate (2). The presence of Cr and of Mo either in the steel or in the  
41 solution has a significant effect on the critical supersaturation and growth of siderite and on  
42 the precipitation kinetics of the colloidal iron carbonate (4, 7, 8). Within this framework,

43 understanding the kinetics of formation and growth of the amorphous carbonate layer that  
44 is presumed to control the dissolution rate comes into focus. The presence of this film is  
45 deduced from the effect on the current (3, 9). A completely passivating film, deduced to be  
46 a 2-layer film comprising hydrous Fe(II,III) oxides, can be formed at room temperature, pH  
47 9.2 and has been studied ellipsometrically (10). However, this is unlikely to be a model for  
48 the dissolving layer formed at high temperature and lower pH. In earlier work, we  
49 presented a grazing-incidence small-angle X-ray scattering (GISAXS) study of the first stages  
50 of siderite crystallisation during CO<sub>2</sub> corrosion of carbon steel (4). We interpreted the  
51 results as indicating three phases in the development of the surface layer: the initial  
52 formation of a discontinuous layer comprising 'blobs' of amorphous FeCO<sub>3</sub>, which  
53 subsequently grow to spread over the surface and eventually cover it; then the appearance  
54 of a new set of scatterers with lateral size scale  $\approx 20$  nm, interpreted as pre-nucleation  
55 clusters; then finally the nucleation and growth of crystals, which is signalled by the  
56 appearance of a crystalline diffraction pattern. The effect of the presence of trace Cr<sup>3+</sup> in  
57 the solution could be interpreted as a significant acceleration and overlapping of all three  
58 stages. The present work reports extension of the GISAXS study to cover different alloys  
59 and develops more detail concerning the initial stages of the corrosion reaction.

## 60 **Experimental**

61 GISAXS experiments were performed on the SAXS/WAXS beamline at the Australian  
62 Synchrotron (11). Steel rods 2 mm in diameter were coated in epoxy and the top surfaces  
63 polished using 1 micron diamond paste to an optically smooth finish. Two types of steel  
64 were studied. The AISI 1006 simple carbon steel (mild steel) sample had an equiaxed  
65 ferritic/pearlitic microstructure in both transverse and longitudinal cross section to the

66 drawing direction. It was relatively free of inclusions and had a grain size of  $\sim 10\ \mu\text{m}$ , typical  
67 of hot formed low carbon steel. The as-received API 5CT L80 1Cr0.25Mo pipeline steel  
68 (supplied by Tenaris) had a quenched and tempered martensitic microstructure, with high  
69 yield strength and mechanical properties appropriate for use for well applications.

70 Micrographs are given in the Supporting Information (SI): figure S1. Each rod was mounted  
71 in an electrochemical cell as used previously for *in situ* synchrotron X-ray diffraction  
72 experiments (6, 12). It featured a solution reservoir, that fitted snugly into a heating jacket,  
73 and which tapered to a narrow slot just above the electrode mounted in the base. The  
74 solution path length at the electrode was approximately 2.5 mm. X-rays of wavelength  
75  $0.8265\ \text{\AA}$  (15 keV) struck the rod at grazing incidence; the exact incidence angle was  
76 unknown, but was determined from the initial position of the diffuse reflected beam to be  
77 less than  $0.5^\circ$ . The synchrotron beam was 0.1 mm wide and spilled over the entire length of  
78 the sample. Scattered X-rays were detected using a Pilatus 1 M detector located 6.73 m  
79 from the sample. This gave a horizontal angular range of  $-0.63 < \theta_f < 0.80$  degrees ( $q$  range  
80 of  $-0.084\ \text{\AA}^{-1} < q_{xy} < 0.107\ \text{\AA}^{-1}$ ) and vertical angular range of  $-0.39 < \alpha_f < 1.14$  degrees ( $q$ -  
81 range of  $-0.051\ \text{\AA}^{-1} < q_z < 0.152\ \text{\AA}^{-1}$ ). The exposure time for each image was 4 s. Thermal  
82 heating by the synchrotron beam is negligible; although the flux is high, the energy  
83 bandwidth and therefore the total power of the beam is low (for a beam of similar flux and  
84 energy focussed into an area a factor of  $10^7$  smaller a temperature rise of 8 K has been  
85 calculated (13)). A more significant problem is the formation of radical species and  
86 hydrogen bubbles due to radiation damage. In preliminary work with images continuously  
87 recorded, we noted using the wide-angle detector, that a diffraction pattern of chukanovite  
88 ( $\text{Fe}_2(\text{OH})_2\text{CO}_3$ ) appeared much more rapidly than expected (an explanation is the effect of a  
89 local pH change that could be caused by radiolysis). When the radiation dose was reduced

90 by attenuating the flux and reducing the image acquisition frequency then the formation of  
91 crystalline products occurred on timescales commensurate with those observed in  
92 experiments in the absence of X-rays (2). Thus, in the present work, the X-ray flux was  
93 reduced to 25% of the full beam using attenuators , images were recorded continuously for  
94 1 minute at the open-circuit potential, the electrode potential was stepped to the desired  
95 value, images were recorded continuously for 2 minutes, then reduced in frequency to one  
96 image every minute. The electrochemical tests were performed using a 0.5 M NaCl solution  
97 that was heated to 80°C and saturated with CO<sub>2</sub> by continuously bubbling CO<sub>2</sub> gas for at  
98 least 1 hr, with pH adjusted by adding 2 M NaOH (20 mL per litre NaCl solution). This gave a  
99 calculated pH of 6.8 at 80°C (6.3 at room temperature). The hot solution was transferred to  
100 the cell and CO<sub>2</sub> was bubbled continuously throughout the experiment. The cell was aligned  
101 in the X-ray beam whilst the temperature stabilised. The open circuit potential was in the  
102 range -0.71 to -0.73 V Ag|AgCl, 3 M KCl. Data processing was performed using software  
103 developed in-house. The raw data images were rotated about the beam centre so that the  
104 specular reflection was aligned in the vertical direction. Cuts at constant  $q_{xy}$  and  $q_z$  could  
105 then be obtained for each image (although in this manuscript we focus mainly on the cuts at  
106 constant  $q_{xy}$ ). The same cell and procedure was used for grazing-incidence X-ray diffraction  
107 measurements, performed separately on the Powder Diffraction (PD) beamline of the  
108 Australian Synchrotron.

109

## 110 **Results and Discussion**

111 GISAXS results from a typical *in situ* experiment, where the sample was corroded under  
112 potentiostatic control, are shown in Figures 1-2. Additional images are given in the SI.

113 Initially, several features were observed: (a) a strong specular streak, which consists of a  
114 peak with a series of additional oscillations in  $q_z$ , (b) broad diffuse scattering around the  
115 specular reflection, and (c) isotropic scattering at low  $q$  near the beam centre. As the  
116 experiment progressed, the specular peak and its oscillations moved to lower  $q_z$ , as did the  
117 diffuse scattering. The amplitude of the oscillations decreased with time. At longer times the  
118 isotropic scattering increased and eventually dominated the signal.

119 The shift in the position of both the specular peak and the oscillations to lower  $q_z$ , and the  
120 damping of the oscillations over time, were particularly clear for measurements conducted  
121 at the open-circuit potential (Figure 2). The time-scale over which the changes occurred  
122 was significantly shorter for samples that were anodically polarised: Figure 3. The  
123 oscillations were observed both in the absence of the solution (figure S2) and in the  
124 presence of the solution when the steel had been cathodically polarised to remove any air-  
125 formed surface film possibly present (figure S3). The timescales for these changes differed  
126 significantly for the mild steel compared to the 1Cr0.25Mo steel (Figure S4)

127 For a sample corroded anodically in 0.03 M HCl + 0.5 M NaCl electrolyte at room  
128 temperature, a similar but much smaller shift in position of the specular reflection was  
129 observed (Figure S5). The oscillations were also present, but less marked, and their damping  
130 with time was also less marked. In this medium, no surface film was expected based on  
131 electrochemical measurements (14); literature X-ray photoelectron spectroscopy results  
132 indicate any film present to have a thickness less than 2-5 nm (15). Thus the differences in  
133 behaviour between this sample and the one corroded in the CO<sub>2</sub> solution are likely due to  
134 the formation of a surface film. However, as we will show, this process involves multiple

135 factors that all contribute to changes in the surface roughness and hence the scattering  
136 features observed.

137 The features we are most interested in are the specular and diffuse scattering. The isotropic  
138 scattering arises from the solution and the cell windows, plus particles suspended in  
139 solution and/or that adhere to the windows of the cell; the formation of these particles  
140 explains why the signal at low  $q$  increases at later times. Features commonly observed in  
141 GISAXS patterns include specular reflections, Yoneda peaks, Kiessig fringes and off-axis  
142 scattering. Specular scattering occurs when the reflection condition is satisfied, i.e. when  
143 the exit angle  $\alpha_f$  is equal to the incident angle  $\alpha_i$ . In most GISAXS studies the incident angle is  
144 chosen to be close to or below the critical angle  $\alpha_c$  for total external reflection of the X-rays,  
145 so that substrate effects are minimised. Yoneda peaks (or 'wings') are observed at an exit  
146 angle  $\alpha_f$  equal to the critical angle  $\alpha_c$ . At this position there is an intensity enhancement due  
147 to the propagation of the evanescent wave. Where there is a thin film on top of a substrate  
148 and the different materials have different indices of refraction for X-rays (and therefore  
149 different critical angles), it is possible to observe multiple Yoneda peaks (16). Kiessig fringes  
150 are oscillations arising from interference from the two interfaces of a thin film (i.e.  
151 substrate-film and film-air). They may be observed either on the specular features (i.e.  
152  $q_{xy}=0$ ) or form factor scattering features (17, 18) . Finally, off-axis scattering arises from the  
153 presence of scattering objects on or embedded in the surface, including pits. This can  
154 comprise both form factor scattering, from the size and shape of the objects, and structure  
155 factor scattering, from their spacing and regularity (19-21).

156 In our scattering data, the specular peak is much broader than expected for a typical GISAXS  
157 pattern. This is due to the fact that the electrodes were not perfectly flat, but, as a result of

158 the polishing method, they were slightly rounded. This means that the angle of incidence  
159 was not well-defined, thus broadening the specular scattering in both the vertical and  
160 horizontal directions. Another contributing factor to the broadness of the specular peak and  
161 the presence of diffuse scatter is the fact that although the electrode was polished to an  
162 optically smooth finish, it was not *atomically* smooth. To test whether random surface  
163 roughness alone could account for the diffuse scatter observed, we used the method of  
164 Ward (22), who gave simple equations in angular space for calculating the diffuse reflection  
165 of light. The detail is given in the SI. The calculation reproduced the observed initial diffuse  
166 scattering, a spot that is broader in the vertical direction than the horizontal, with small  
167 values for the RMS roughness (SI). The broadening of the specular reflection also means it is  
168 unlikely that we would be able to observe any Yoneda features in our scattering data, as  
169 these will also be broadened and/or obscured.

170 The presence of oscillations along the specular direction is reminiscent of Kiessig fringes  
171 arising from interference by a thin film, however, they were also observed when no surface  
172 film was present: (a) the as-prepared electrodes with no solution present (Fig. S2); (b) at the  
173 start of an experiment, before any corrosion product nuclei would have been formed on the  
174 surface (Figure S3); and (c) while the steel electrode was negatively polarised (even to the  
175 extent of significant hydrogen evolution) to remove any air-formed surface film that might  
176 possibly be present (Fig. S3). Given that for a material such as our steel, having a complex  
177 microstructure, some elements of which have different X-ray scattering characteristics, a  
178 plausible explanation is that the oscillations are due to particles of  $\text{Fe}_3\text{C}$  emerging at the  
179 surface, i.e. at the grain boundaries, the size-scale of the steel microstructure being  
180 commensurate with the scale probed by the GISAXS experiment. The calculated scattering  
181 from such second-phase particles (Figure S6) exhibits form factor oscillations along the



182 specular peak that are similar to the experimental data when superimposed on the broad  
183 specular scattering peak. The detail of the calculation is given in the SI, using X-ray scattering  
184 parameters appropriate for Fe and Fe<sub>3</sub>C. Thus, the oscillations cannot be taken as evidence  
185 for a film on the electrode surface.

186 We therefore attribute the scattering features as follows: The broad specular peak and the  
187 diffuse scattering arise from the surface roundedness and roughness, while the oscillations  
188 arise from the steel microstructure (the distribution of Fe<sub>3</sub>C). However, neither of these  
189 explanations can account for the movement of the specular peak (and the oscillations) to  
190 lower angles as the experiment progressed. In fact, the only explanation for the movement  
191 of the specular peak is for the incidence angle to change. Another observation to note is the  
192 damping of the oscillations as the experiment progressed.

193 In the initial stages of the experiment, the metal dissolves, assumed through the formation  
194 of a thin, amorphous surface film (2). The surface roughness would increase. Two factors  
195 could contribute to an increase of surface roughness: the formation of a surface deposit  
196 that is itself rough; and smaller metal crystals dissolving in preference to larger ones,  
197 perhaps as a consequence of the dissolution rate constant varying with position away from  
198 a grain boundary or inclusion, or different metal crystal faces dissolving at different rates  
199 (23). In addition, the Fe<sub>3</sub>C particles do not dissolve. As a result, the corroded microstructure  
200 may have sufficiently deep features to cause a change in the average angle at which the  
201 incident beam strikes the surface. If these surface features were much larger than the X-ray  
202 wavelength, then a classical approach can be used to calculate the scattering, based on  
203 geometry and the law of reflection. We simply modelled the developing surface roughness  
204 as a sine function, superimposed on a second sine function to model the slight rounding of

205 the surface due to the polishing, and considering the initial macroscopic roughness just as  
206 something which would impose a diffuse scattering on top of the calculated effect. The  
207 detail is in the SI. The calculated effect of a sinusoidal surface roughening, now on the nm  
208 scale, is to move the peak position to lower angles. This occurs both when the number of  
209 periods and when the amplitude increase. Thus surface roughening can qualitatively explain  
210 the experimentally observed movement of the peak to lower  $q_z$  as the experiment  
211 progressed. Given the sine wave approximation for the surface profile, the change in  
212 reflection angle from the start of the experiment can be related to the change in average  
213 gradient of the surface with respect to the surface plane, which is the chord of the sine  
214 function,  $\Delta g = \Delta(2h/\Lambda)$ , where  $h$  denotes the amplitude and  $\Lambda$  the wavelength of the  
215 model sine function. Figure 4 compares the results for the two steels, showing also the  
216 variation of the current with time, and the variation in the fractal dimension derived from  
217 the development of 'spottiness' in the diffraction ring for Fe, measured in a separate  
218 experiment (24). The fractal dimension is related to the distribution of crystallite sizes. An  
219 increase in the fractal dimension corresponds to an increase in the average crystallite size  
220 probed by the X-ray beam. In our earlier work, this was attributed to preferential dissolution  
221 of smaller Fe grains (6, 23, 24). Figure 5a shows the variation of  $\Delta g$  with time, comparing the  
222 two different steels under anodic polarisation and also comparing variation for the mild  
223 steel at the open circuit condition with that under anodic polarisation. Figure 5b compares  
224 the variation of  $\Delta g$  for the two steels under CO<sub>2</sub> corrosion with the mild steel under anodic  
225 corrosion in acidic NaCl solution (scattering data in Figure S5).

226 The results show a semi-logarithmic variation of the surface roughness measure,  $\Delta g$ , over  
227 the whole time when the current was constant, before a detectable crystalline film

228 nucleated on the surface. The consistency of electrochemical results taken on the two  
229 different synchrotron visits is notable (current densities shown in Figure 4). The  
230 semilogarithmic time variation illustrates the very rapid development of very small size-  
231 scale features, which then slows with time. The variation in reflection angle in the GISAXS  
232 experiment reveals the development of surface roughness on the nm scale. That deduction  
233 is based on the assumption of a size scale for  $\lambda$ . If we follow through an assumption that  
234 the small-scale roughness is caused by the growth of a surface film, then it is reasonable to  
235 assume that the growth might be modulated by crystallite grain boundaries emerging at the  
236 surface – because growth or dissolution can be different on different crystal faces or at grain  
237 boundaries (25-29)- or by the emergence of Fe<sub>3</sub>C at the surface. If  $\lambda$  is taken as the average  
238 crystallite size or distance between Fe<sub>3</sub>C plates revealed in the micrographs of the etched  
239 surfaces (~10  $\mu$ m) then  $h$  is of order 10nm, since  $\Delta g$  is of order  $10^{-3}$ . However, the result can  
240 also be explained with much smaller scales of  $\lambda$  and  $h$ : the simulations given in the SI show  
241 that the observed variation of  $\Delta g$  could also be obtained with  $h$  of scale 1 nm or less, that is  
242 with  $\lambda$  of scale 0.1 – 1  $\mu$ m. The growth of anodic oxides results in a roughness developing  
243 on the nm scale, even on single-crystal substrates (30) so these size scales for the  
244 developing surface roughness are not unexpected for formation of an anodic layer. Figure  
245 5a shows that the time scale for the development of surface roughness for the mild steel  
246 decreases with increasing anodic polarisation, as would be expected for a surface roughness  
247 change driven by an anodic reaction.

248 The comparison between the average surface gradient developed under CO<sub>2</sub> corrosion and  
249 that developed under anodic dissolution in acidic solution (Figure 5b) in part discriminates  
250 between roughness developed as a consequence of anodic dissolution (e.g. different

251 dissolution rates of different size crystallites, or different dissolution rates close to grain  
252 boundaries or inclusions) from roughness developed as a consequence of surface film  
253 growth. The acid dissolution was performed under galvanostatic conditions, at first at a low  
254 current density, to avoid any precipitation of  $\text{FeCl}_2$  and then at a similar total dissolution  
255 rate to that found in the  $\text{CO}_2$  experiments, and the comparison is made in terms of total  
256 charge passed, again to compare surface gradient developed with the same amount of  
257 material dissolved. The measured changes were much smaller, implying that the initial  
258 development of surface roughness in the  $\text{CO}_2$  corrosion experiment (on a nm scale) can be  
259 attributed to the development of a thin but rough surface layer.

260 We have extensively discussed the nucleation and growth of the crystalline siderite,  
261 signalled by the current peak for the mild steel (2, 4, 6-8, 12, 24, 31-33). On the 1Cr0.25Mo  
262 pipeline steel the current peak is absent and we have interpreted this in terms of the effect  
263 of Cr and Mo decreasing the critical supersaturation for formation of crystalline siderite (4,  
264 7, 8). Figures 4 and 5 show that, although the total current density for these two steels was  
265 very similar, the time scale for development of surface roughness of the 1Cr0.25Mo pipeline  
266 steel was shorter than that for the mild steel, at the same electrode potential, though the  
267 slope of the semilogarithmic time dependence is similar. It would appear that the very  
268 smallest roughness scales appear more rapidly on the 1Cr0.25Mo pipeline steel than on the  
269 mild steel. One hypothesis, which we have previously advanced, is that these smallest  
270 roughness scales reflect the nucleation and initial growth of amorphous ferrous carbonate,  
271 in the form of nm-scale lumps (4), similar to models advanced for the initial stages of anodic  
272 film formation by vertical growth and lateral spreading from nucleation centres (34). The  
273 more rapid appearance of roughness on this size scale on the 1Cr0.25Mo pipeline steel than  
274 on the mild steel is consistent with the hypothesis that Cr and Mo in the steel accelerate the

275 precipitation of amorphous iron carbonate and increase the growth rate, eventually leading  
276 to microcrystalline protective layers.

277 The damping of the oscillations in the scattering pattern in the very first stages of  
278 dissolution, as dissolution progresses, implies that the microstructural contrast between  
279 Fe<sub>3</sub>C and Fe in the substrate steel that gives rise to these oscillations becomes obscured. The  
280 simplest explanation is that the X-rays reflected from the substrate become attenuated in  
281 passing through a surface layer. The surface layer could be simply due to the physical  
282 roughness – a mixture of Fe and water - or could be the amorphous FeCO<sub>3</sub> that is assumed  
283 to be present, if this were of sufficient thickness. At this X-ray wavelength, the attenuation  
284 length of either FeCO<sub>3</sub> or of a layer that is 50 vol% Fe with water is of order 100 μm. The  
285 angle of incidence of X-rays on the surface is in the range  $\alpha_i = 0.1 - 0.5^\circ$ . Thus the path  
286 length of X-rays through a surface layer of thickness  $s$  would be  $2s / \sin \alpha_i$ . Given an  
287 attenuation length of 100 μm, damping of the reflection oscillations would imply  $s \sim 0.1 -$   
288  $0.5 \mu\text{m}$ . This is larger though still consistent with the presumed size scale for development  
289 of the surface roughness, discussed above, within the limitations of the approximations  
290 used to interpret the data.

291 Whereas GISAXS reveals the variation in morphological roughness, the variation in fractal  
292 dimension of the Fe diffraction peak revealed by the grazing-incidence diffraction  
293 experiment, that develops later in the course of the experiment, was interpreted in terms of  
294 a change in crystallite size distribution intersecting the X-ray beam: that is, a corrosion-  
295 induced surface roughening by removal of material from the surface, which is distinct from  
296 the model of surface film growth advanced for the interpretation of the GISAXS data.  
297 Interpretation of the variation of fractal dimension required a crystallite size scale of

298 dissolving material in the range 0.1 to 10  $\mu\text{m}$ , of relative amount  $10^4: 10^3: 1$  for 0.1, 1 and 10  
299  $\mu\text{m}$  features respectively at the start changing to  $10^3: 10^2: 1$  at the end (23). The spatial  
300 heterogeneity of dissolution rate, on size scales smaller than are revealed in electron  
301 micrograph images of the etched microstructure, is remarkable. The semilogarithmic  
302 variation with time can be interpreted as reflecting the frequency distribution of the metal  
303 crystallites of different size scale, with the progressive dissolution of the smaller crystallites  
304 leaving the surface roughness dominated by the larger ones. This change in crystallite size  
305 distribution at the surface (which is presumed to relate to the morphological roughness)  
306 develops significantly more slowly on the 1Cr0.25Mo steel than on the mild steel, because  
307 siderite crystals nucleate and grow more rapidly, causing the current to fall as the  
308 dissolution becomes blocked by the crystalline layer, as previously shown.

309 This work has illustrated the difficulties and subtleties of interpretation that appear when  
310 GISAXS is applied to *in situ* electrochemical studies of polycrystalline materials. It has also  
311 illustrated that appropriate comparison can reveal significant detail concerning the  
312 development of surface layers on nm thickness scales, even when the system is not ideal.  
313 Previously presented interpretations of the effects of microalloying on the formation of  
314 surface layers during  $\text{CO}_2$  corrosion of steel in hot brine have been strengthened. The  
315 method has some application in evaluating the effects of different microalloying additions  
316 to the steel, particularly in following the initial slow evolution of the surface at the corrosion  
317 potential, before the nucleation and growth of the protective crystalline siderite.

318

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333 **Competing Interests:** The authors declare no competing interests

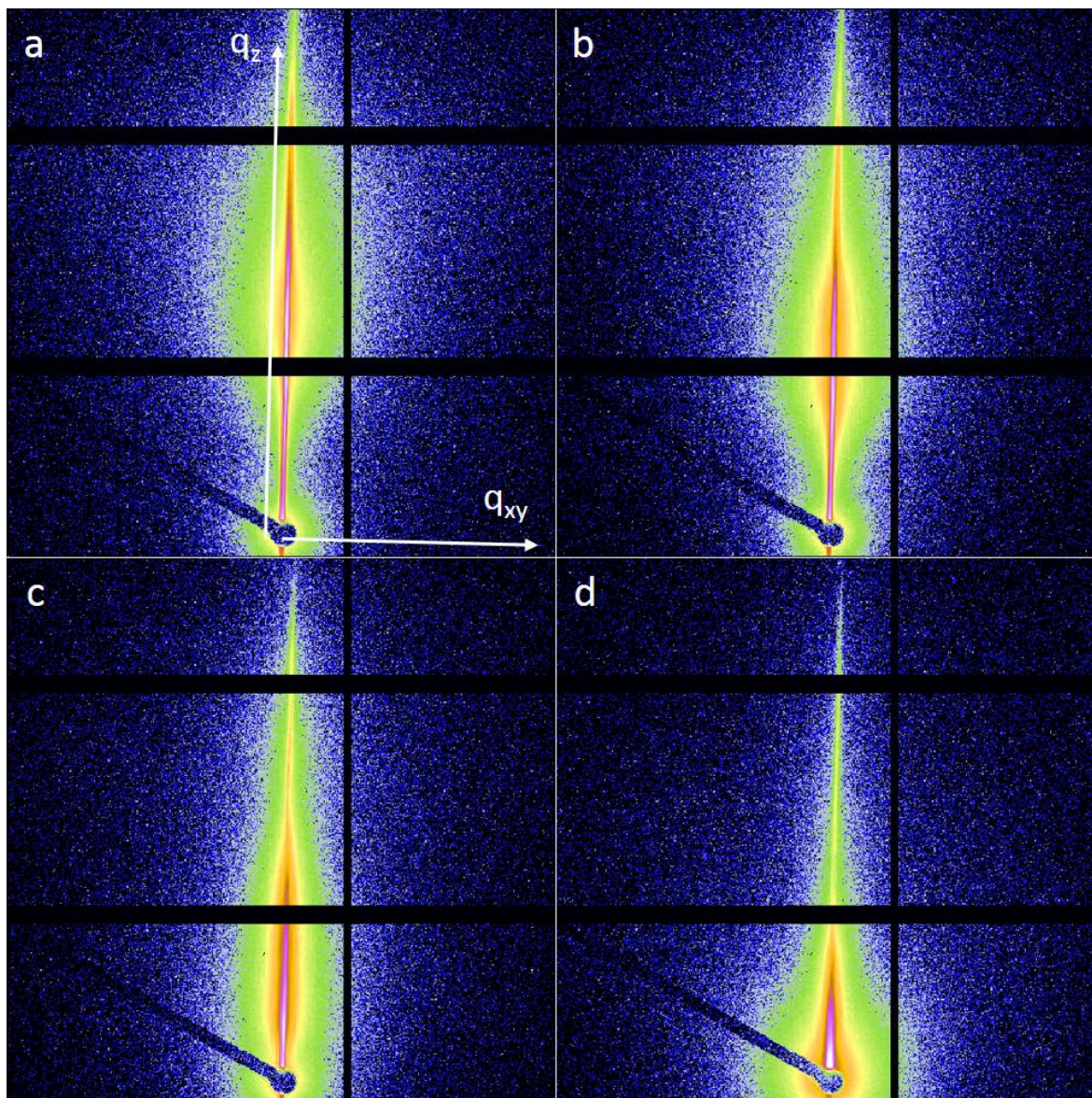
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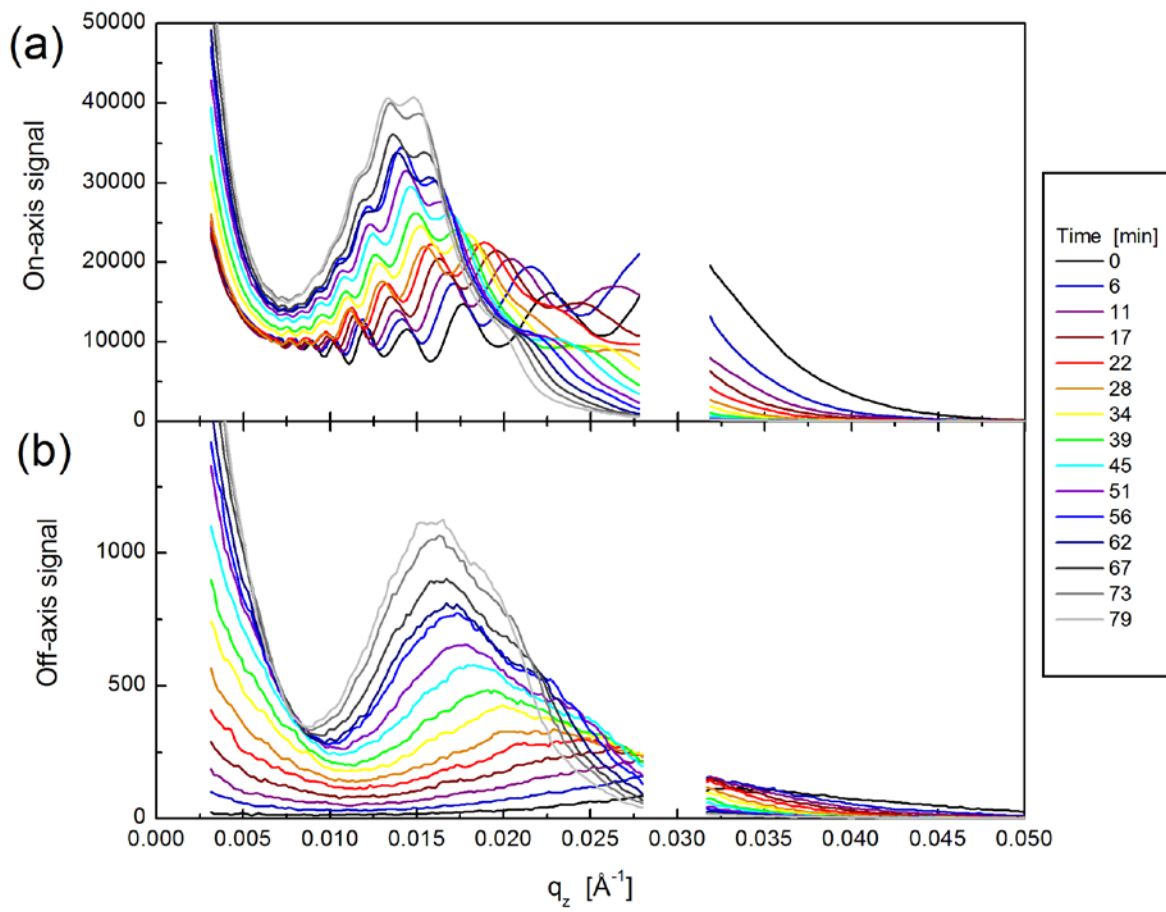




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388 Figure 1. Selected raw GISAXS images from an experiment where a mild steel electrode was  
 389 corroded electrochemically at o.c.p. + 200 mV in a 0.5 M NaCl, CO<sub>2</sub>-saturated solution with  
 390 pH = 6.8 at 80°C, at times from application of current: (a) 0 s, (b) 18 s, (c) 110 s, (d) 600 s  
 391 (o.c.p denotes open circuit potential). The  $q_{xy}$  and  $q_z$  directions are indicated in (a). The  
 392 horizontal and vertical black bars are gaps between the detector modules, while the  
 393 diagonal bar is the beam stop.

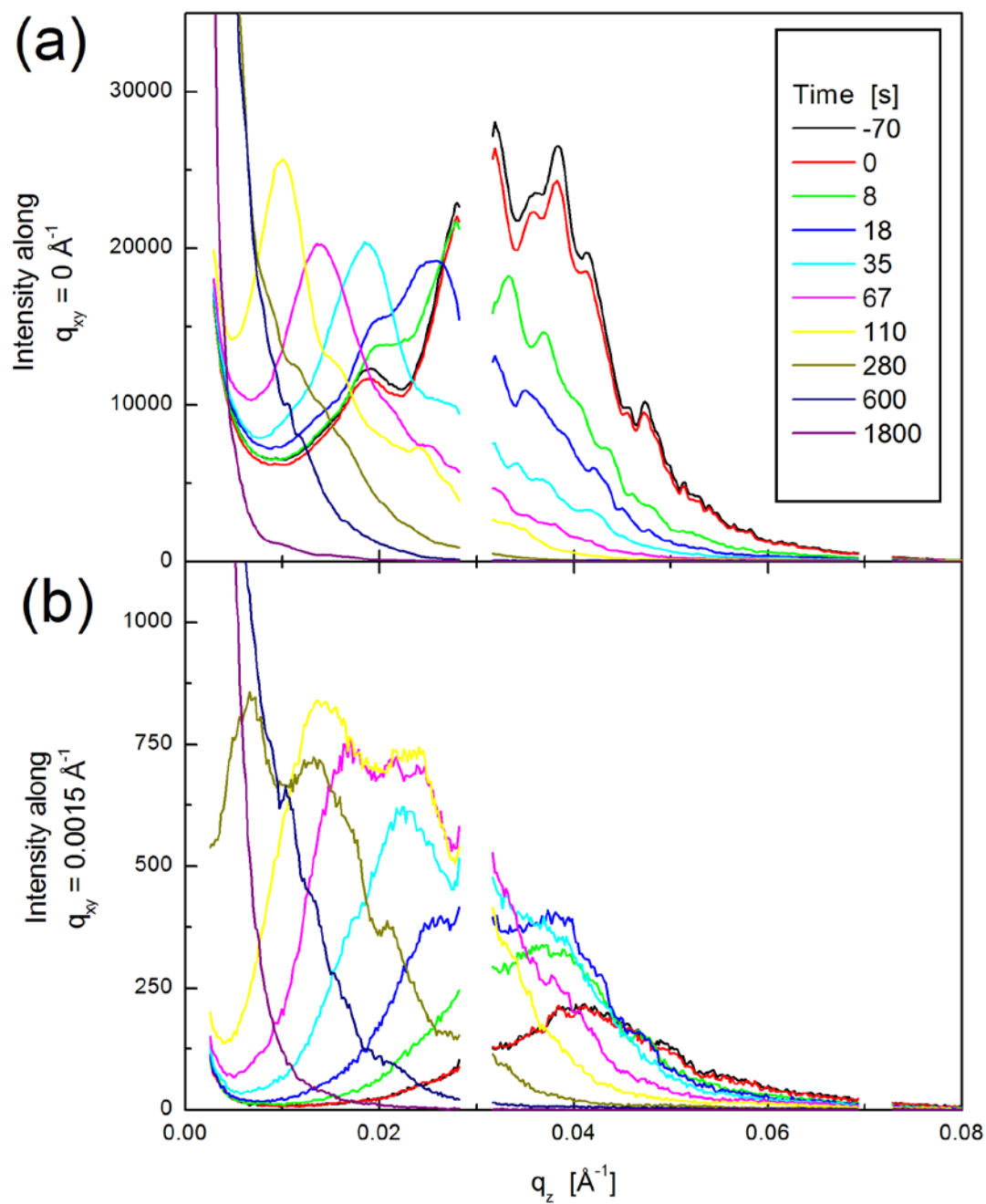
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397 Figure 2. GISAXS intensity versus  $q_z$ , for mild steel corroding at the o.c.p. (-720 mV vs.  
 398 Ag/AgCl) in a 0.5 M NaCl, CO<sub>2</sub>-saturated solution with pH = 6.8 at 80°C. Two cuts at constant  
 399  $q_{xy}$  for a selection of scans is shown: (a) the specular (on-axis) scattering at  $q_{xy} = 0 \text{ \AA}^{-1}$  and (b)  
 400 the diffuse (off-axis) scattering at  $q_{xy} = 0.002 \text{ \AA}^{-1}$ .

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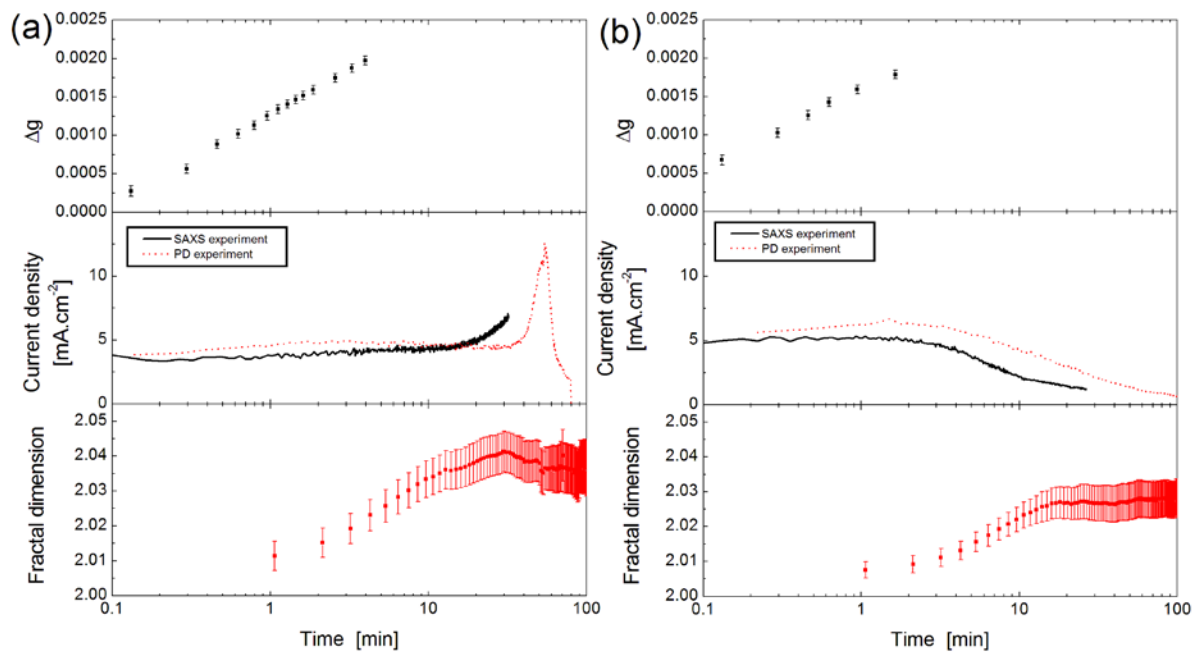


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404 Figure 3. GISAXS intensity versus  $q_z$ , extracted from (a) the specular scattering,  $q_{xy} = 0 \text{ \AA}^{-1}$ ,405 and (b) the diffuse scattering,  $q_{xy} = 0.0015 \text{ \AA}^{-1}$ , for selected scans from the experiment

406 shown in Figure 1, at times as indicated.

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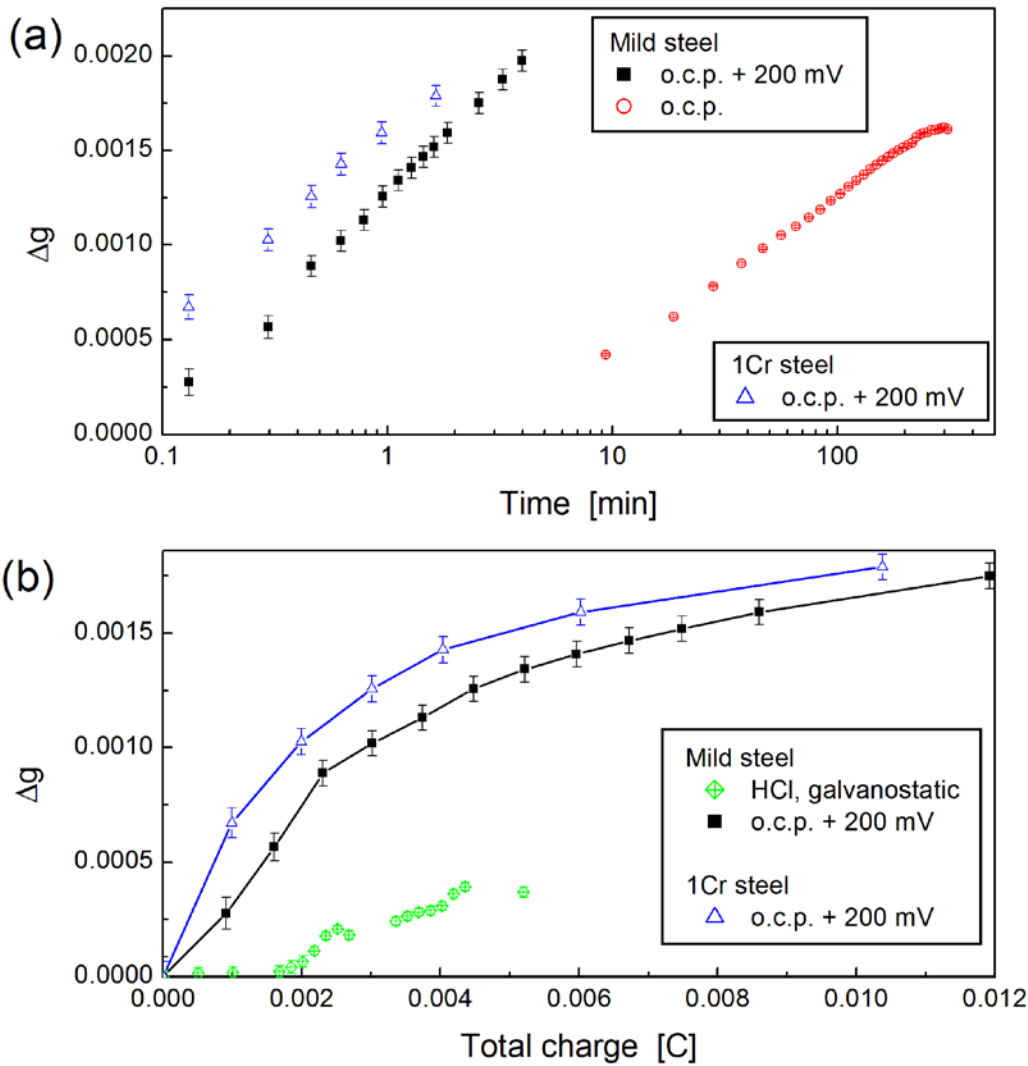


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410 Figure 4. Variation with time of average surface gradient (for data where the specular peak  
 411 could be resolved), current density and fractal dimension, for (a) AISI 1006 mild steel and (b)  
 412 API 5CT L80 1Cr0.25Mo pipeline steel under anodic polarisation in CO<sub>2</sub>-saturated 0.5M NaCl  
 413 at 80°C. The results are from separate synchrotron experiments carried out several months  
 414 apart.

415



416

417 Figure 5. (a) Variation with time of average surface gradient, for AISI 1006 mild steel and API

418 5CT L80 1Cr0.25Mo pipeline steel, at the open circuit corrosion potential and under anodic

419 polarisation in  $\text{CO}_2$ -saturated 0.5M NaCl at 80°C. (b) Variation with total anodic charge

420 passed of average surface gradient, for AISI 1006 mild steel and API 5CT L80 1Cr0.25Mo

421 pipeline steel, under potentiostatic anodic polarisation in  $\text{CO}_2$ -saturated 0.5M NaCl at 80°C

422 and for AISI 1006 mild steel under galvanostatic polarisation at similar current density in

423 0.03 M HCl + 0.5 M NaCl electrolyte at room temperature.

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